

# Wavelet Analysis and Filtering to Identify

# Dominant Orientations of Permeability

## Anisotropy<sup>1</sup>

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### Abstract

An accurate representation of permeability anisotropy is needed to correctly model the rate and direction of groundwater flow. We develop a wavelet analysis technique that can be used to characterize principal directions of anisotropy in both stationary and non-stationary permeability fields. Wavelet analysis involves the integral transform of a field using a wavelet as a kernel. The wavelet is shifted, scaled, and rotated to analyze different locations, sizes, and orientations of the field. The wavelet variance is used to identify scales and orientations that are dominant anywhere in the field. If the field is non-stationary, such that different zones of the field are characterized by different dominant scales or orientations, the wavelet variance can identify all dominant scales and orientations if they are distinct. If the dominant scales and orientations of different zones are similar, the wavelet variance only identifies the dominant scale and orientation of the primary zone. In this paper, we present a combined wavelet analysis and filtering approach to identify all dominant scales and orientations in a non-stationary permeability field. We apply the method to laboratory-collected permeability data from Massillon sandstone.

*Key words:* heterogeneity, non-stationary, Morlet wavelet

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23 Porous medium properties, such as permeability, are often spatially variable  
 24 and anisotropic. In a layered porous medium, permeability varies with direc-  
 25 tion, with the highest and lowest permeabilities found in the directions paral-  
 26 lel and perpendicular to layering, respectively. These directions are called the  
 27 principal directions of anisotropy.

28 Fluid flow in porous media is described by Darcy's law, given by

$$29 \begin{bmatrix} q_x \\ q_y \end{bmatrix} = -\frac{\rho g}{\mu} \begin{bmatrix} k_{xx} & k_{xy} \\ k_{yx} & k_{yy} \end{bmatrix} \begin{bmatrix} \partial h / \partial x \\ \partial h / \partial y \end{bmatrix}, \quad (1)$$

30 where  $q_x$  and  $q_y$  are components of specific discharge in the  $x$  and  $y$  directions,  
 31 respectively,  $\rho$  is the fluid density,  $g$  is the gravitational constant,  $\mu$  is the  
 32 fluid viscosity,  $k_{ij}$  is the  $i,j$ th entry of the permeability tensor, and  $\partial h / \partial x$  and  
 33  $\partial h / \partial y$  are the hydraulic gradients in the  $x$  and  $y$  directions, respectively. The  
 34 components of the permeability tensor are given by (Bear, 1972)

$$35 \quad k_{xx} = \frac{k_{\parallel} + k_{\perp}}{2} + \frac{k_{\parallel} - k_{\perp}}{2} \cos(2\theta) \quad (2)$$

$$36 \quad k_{yy} = \frac{k_{\parallel} + k_{\perp}}{2} - \frac{k_{\parallel} - k_{\perp}}{2} \cos(2\theta) \quad (3)$$

$$37 \quad k_{xy} = k_{yx} = \frac{k_{\parallel} - k_{\perp}}{2} \sin(2\theta) \quad (4)$$

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40 where  $k_{\parallel}$  is the average permeability parallel to layering,  $k_{\perp}$  is the average  
41 permeability perpendicular to layering, and  $\theta$  is the orientation of layering,  
42 which we call the dominant orientation. If the assumed value of the dominant  
43 orientation of the permeability tensor is incorrect, the magnitude and direc-  
44 tion of the specific discharge from (1) will also be incorrect, with errors in  
45 magnitude of up to 30% and errors in flow direction of up to  $45^{\circ}$  (Anderman  
46 et al., 2002). The errors increase as the anisotropy ratio,  $k_{\parallel}/k_{\perp}$ , increases.

47 Several studies demonstrate the importance of an accurate representation of  
48 aquifer anisotropy. Anisotropy was shown to have significant effects on the  
49 patterns of groundwater seepage from lakes (Genereux and Bandopadhyay,  
50 2000), on groundwater travel times in sedimentary fractured rocks (Burton et  
51 al., 2002), on seepage in bog peat (Beckwith, Baird, and Heathwaite, 2003),  
52 and on the migration processes of infiltrated stream water to a partially pen-  
53 etrating well (Chen and Chen, 2003).

54 In this paper, we present a method for identifying dominant orientations in  
55 a non-stationary permeability field that exhibits anisotropy. We consider the  
56 class of non-stationarity in which the permeability field contains zones with  
57 different dominant orientations and scales, but within a given zone, the prop-  
58 erties are stationary. We assume that point measurements of permeability are  
59 isotropic, but that adjacent permeability measurements are correlated with  
60 different correlation lengths in different directions. Thus, the permeability field  
61 exhibits structural anisotropy.

62 Two methods that have been used to determine dominant orientations are di-  
63 rectional variograms and anisotropic wavelet analysis. A directional variogram  
64 identifies spatial correlation by estimating the variability in permeability as a

65 function of the separation distance between measurements and of the relative  
66 orientation of the separations. The dominant orientation is the direction that  
67 exhibits correlation through the longest separation distance (Isaaks and Sri-  
68 vastava, 1989). Tidwell and Wilson (2000) use variograms to quantify spatial  
69 variability in permeability measurements from a block of Massillon sandstone.  
70 The sandstone exhibits layering of bounding surfaces with low permeability  
71 that separate cross-stratification sets that have higher permeability. A two-  
72 dimensional variogram identified the dominant orientations of each rock face.

73 Neupauer et al. (2006) developed a wavelet analysis approach to identify dom-  
74 inant orientations of an anisotropic permeability field. In anisotropic wavelet  
75 analysis, an integral transform is performed on the permeability data using an  
76 anisotropic kernel. The largest value of the integral transform occurs when the  
77 orientation of the kernel matches the dominant orientations of the permeabil-  
78 ity field. If the permeability field is non-stationary, with different dominant  
79 orientations at different locations, the method of Neupauer et al. (2006) iden-  
80 tified all dominant orientations only if the dominant orientations are distinct.  
81 If the dominant orientations are not distinct, the method only identified the  
82 primary (most dominant) orientation, while secondary orientations could not  
83 be identified.

84 In this paper, we enhance the wavelet analysis method of Neupauer et al.  
85 (2006) to develop a method that identifies both primary and secondary ori-  
86 entation in a non-stationary, anisotropic permeability field. This method uses  
87 a combination of wavelet analysis and filtering. The method has been de-  
88 veloped for two-dimensional, uniformly-spaced permeability data sets. In the  
89 next section, we present wavelet analysis theory, and we introduce the filtering  
90 procedure. In the subsequent section, we illustrate the new combined wavelet

91 analysis and filtering method using two-dimensional synthetic data sets. Fi-  
92 nally, we apply the method to Tidwell and Wilson’s (2000) permeability data  
93 from Massillon sandstone.

## 94 METHODOLOGY

95 The methodology used to implement wavelet analysis on a two-dimensional  
96 field is explained in this section. We present general wavelet analysis theory,  
97 followed by our specific procedure for identifying dominant orientations in a  
98 non-stationary random field.

### 99 Theory

100 The continuous wavelet transform of a two-dimensional field,  $f(\mathbf{x})$ , is given by  
101 (e.g., Farge, 1992)

$$102 \quad W^\theta(a, \mathbf{b}, \theta; L) = \int f(\mathbf{x}) \psi_{a, \mathbf{b}, \theta; L}^*(\mathbf{x}) d\mathbf{x}, \quad (5)$$

103 where  $W^\theta(a, \mathbf{b}, \theta; L)$  is the wavelet coefficient,  $\psi_{a, \mathbf{b}, \theta; L}(\mathbf{x})$  is the scaled, shifted,  
104 and rotated two-dimensional wavelet,  $a$  is the scaling factor,  $\mathbf{x} = (x, y)$  is the  
105 spatial domain,  $\mathbf{b}$  is the shift vector on the spatial domain  $(x, y)$ ,  $\theta$  is the  
106 angle of orientation of the wavelet relative to the  $+x$  axis (positive  $\theta$  is in the  
107 counterclockwise direction),  $L$  is the anisotropy ratio, defined as the ratio of  
108 the scaling factor in the direction perpendicular to  $\theta$  to the scaling factor in  
109 the  $\theta$  direction, and the superscript asterisk denotes the complex conjugate.  
110 The limits of integration are  $-\infty$  to  $+\infty$  unless otherwise stated. The scaled,

111 shifted, and rotated wavelet is given by (Farge, 1992)

$$112 \quad \psi_{a,\mathbf{b},\theta;L}(\mathbf{x}) = \sqrt{\det(\mathbf{A})}\psi(\mathbf{A}\mathbf{C}(\mathbf{x} - \mathbf{b})), \quad (6)$$

113 where  $\psi(\mathbf{x})$  is the mother wavelet,  $\det(\mathbf{A})$  is the determinant of the matrix  $\mathbf{A}$ ,

114 and  $\mathbf{A}$  and  $\mathbf{C}$  are anisotropy and linear transformation matrices, respectively,

115 given by

$$116 \quad \mathbf{A} = \frac{1}{a} \begin{bmatrix} L & 0 \\ 0 & 1 \end{bmatrix}, \quad (7)$$

117 and

$$118 \quad \mathbf{C} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}. \quad (8)$$

119 A wavelet is a function that has unit energy ( $\int |\psi(\mathbf{x})|^2 d\mathbf{x} = 1$ ), a zero mean,

120 and is non-zero over a finite region (Farge, 1992). In this work we use the

121 Morlet wavelet, defined as (e.g. Farge, 1992)

$$122 \quad \psi(\mathbf{x}) = \frac{1}{\sqrt{\pi}} e^{i\mathbf{k}_o \cdot \mathbf{x}} e^{-\frac{1}{2}(\mathbf{x} \cdot \mathbf{x})}, \quad (9)$$

123 where  $\mathbf{k}_o = [0, k_o]$ , and  $k_o > 5.5$ . We use  $k_o = 2\pi$ . Figure 1 shows the real part

124 of the Morlet wavelet with the effects of changing the scale ( $a$ ), orientation

125 ( $\theta$ ), and anisotropy ratio ( $L$ ). Because the wavelet is non-zero only over a

126 finite region, the wavelet transform identifies local properties of the field. The

127 integral transform in (5) is evaluated for a range of shift parameters,  $\mathbf{b}$ , a range

128 of scale parameters,  $a$ , and a range of orientations,  $\theta$ . Shifting the wavelet in

129 (5) results in the wavelet analysis of different regions of the field; scaling the

130 wavelet analyzes varying sizes of regions within the field, and rotating the

131 wavelet analyzes different orientations of the field. In our analysis, we hold  $L$

132 constant.

133 Continuous wavelet analysis transforms a two-dimensional field,  $f(\mathbf{x})$ , into a  
134 four-dimensional  $(a, \mathbf{b}, \theta)$  wavelet coefficient. To reduce the dimensionality, the  
135 four-dimensional wavelet coefficient is integrated over the  $\mathbf{b}$  domain to obtain  
136 the wavelet variance,  $WV_f(a, \theta)$ , given by (Antoine et al., 2004)

$$137 \quad WV_f(a, \theta) = \int_{\Omega} |W^\theta(a, \mathbf{b}, \theta; L)|^2 d\mathbf{b}, \quad (10)$$

138 where  $\Omega$  is the domain of  $\mathbf{b}$ . Because integration is carried out over the  $\mathbf{b}$   
139 (spatial) domain, all local information is lost; thus the wavelet variance is  
140 a global measure. Large values of the wavelet variance occur at  $(a, \theta)$  pairs  
141 that correspond to scales and orientations that exist anywhere in the field  
142  $f(\mathbf{x})$ . Let us define these dominant scales and orientations as  $a_{\max}$  and  $\theta_{\max}$ ,  
143 respectively. These are globally-dominant scales and orientations. The locally-  
144 dominant orientation at each position  $\mathbf{b} \in \mathbf{x}$  is the orientation,  $\theta$ , at which  
145 the wavelet coefficient achieves its maximum value for a scale of  $a_{\max}$ , and is  
146 given by

$$147 \quad \theta_{\max}(\mathbf{b}) = \{\theta : W^\theta(a_{\max}, \mathbf{b}, \theta; L) = \max(W^\theta(a_{\max}, \mathbf{b}, \theta; L))\}. \quad (11)$$

## 148 Procedure

149 Neupauer et al. (2006) used wavelet analysis to identify dominant orientations  
150 of permeability anisotropy according to the following procedure:

- 151 (1) Use (5) to calculate the wavelet coefficients for a range of shifts ( $\mathbf{b} \in \mathbf{x}$ ),  
152 scales, and orientations, and for a particular anisotropy ratio,  $L$ .
- 153 (2) Use these wavelet coefficients in (10) to compute the wavelet variance.

154 (3) Determine the globally-dominant scales and orientations  $(a_{\max}, \theta_{\max})$  by  
 155 identifying the  $(a, \theta)$  pairs that correspond to local maxima of the wavelet  
 156 variance.

157 (4) Identify the locally-dominant orientations for each globally dominant  
 158 scale,  $a_{\max}$ , using (11).

159 If a field contains two or more globally-dominant scales and orientations  
 160  $(a_{\max}, \theta_{\max})$ , the wavelet variance should have local maxima at each globally-  
 161 dominant pair. However, if the  $(a_{\max}, \theta_{\max})$  pair is not sufficiently distinct, the  
 162 wavelet variance may only have one local maximum at an  $(a_{\max}, \theta_{\max})$  that  
 163 corresponds to the scale and orientation of the primary feature (Neupauer  
 164 et al., 2006). To address this situation, we have developed a continuation of  
 165 the procedure to identify  $(a_{\max}, \theta_{\max})$  for a secondary feature. The procedure  
 166 continues as follows:

167 (5) If the wavelet variance has only one local maximum, filter out the wavelet  
 168 coefficients that correspond to  $(a_{\max}, \theta_{\max})$  to remove the primary feature  
 169 from the original field.

170 (6) Reconstruct the field with the remaining wavelet coefficients. For anisotropic  
 171 wavelet analysis, the field  $f(\mathbf{x})$  can be reconstructed from its wavelet co-  
 172 efficients using (adapted from Chui, 1992; Farge, 1992)

$$173 \quad f(\mathbf{x}) = \frac{2}{\sqrt{LC_\delta}} \int_0^\pi \int_0^\infty \frac{W^\theta(a, \mathbf{b}, \theta; L)}{a^2} da d\theta, \quad (12)$$

174 where

$$175 \quad C_\delta = \iint \frac{\hat{\psi}(u, v)}{u^2 + v^2} dudv, \quad (13)$$

176 where  $\hat{\psi}(u, v)$  is the Fourier transform of the wavelet, and  $u$  and  $v$  are  
 177 transform variables.

178 (7) Repeat Steps 1–4 for the reconstructed field to identify the secondary  
 179 dominant orientation.

## 180 **Implementation**

181 To perform wavelet analysis, the integral transform in (5) is analyzed for a  
 182 range of scales, shifts, and orientations. To reduce the computational burden,  
 183 the integral transform can be evaluated in Fourier space, which eliminates the  
 184 need to integrate over the range of shifts,  $\mathbf{b}$ . Substituting (6) into (5) gives

$$185 \quad W^\theta(a, \mathbf{b}, \theta; L) = \sqrt{\det(\mathbf{A})} \int f(\mathbf{x}) \psi^*(\mathbf{A}\mathbf{C}(\mathbf{x} - \mathbf{b})) d\mathbf{x}. \quad (14)$$

186 The integral in this equation can be written as a convolution, leading to

$$187 \quad W^\theta(a, \mathbf{b}, \theta; L) = \sqrt{\det(\mathbf{A})} f(\mathbf{b}) * \psi^*(-\mathbf{A}\mathbf{C}\mathbf{b}), \quad (15)$$

188 where  $*$  denotes convolution. The convolution can be evaluated efficiently in  
 189 Fourier space as a product of the Fourier transforms of  $f(\mathbf{b})$  and  $\psi^*(-\mathbf{A}\mathbf{C}\mathbf{b})$ ;  
 190 thus the wavelet transform becomes

$$191 \quad W^\theta(a, \mathbf{b}, \theta; L) = \sqrt{\det(\mathbf{A})} \mathcal{F}^{-1} \{ \mathcal{F}[f(\mathbf{b})] \mathcal{F}[\psi^*(-\mathbf{A}\mathbf{C}\mathbf{b})] \}, \quad (16)$$

192 where  $\mathcal{F}$  and  $\mathcal{F}^{-1}$  denote the Fourier and inverse Fourier transforms, respec-  
 193 tively. This equation is evaluated for a range of scales and orientations.

195 In this section we demonstrate the wavelet analysis procedure using two-  
 196 dimensional sinusoidal fields defined as

$$197 \quad z_i(x, y) = \sin \left[ \frac{2\pi}{T_i} (x \cos \phi_i + y \sin \phi_i) \right] \quad (17)$$

198 where  $\phi_i$  is the orientation angle (defined as positive in the counterclockwise  
 199 direction, with  $\phi_i = 0$  in the direction of  $+x$ ),  $T_i$  is the period,  $i$  denotes the  
 200 case number, and the domain is  $-5 \text{ m} < x < 5 \text{ m}$  and  $-5 \text{ m} < y < 5 \text{ m}$ .  
 201 We consider two cases. In Case 1 (Fig. 2A), the field contains two zones with  
 202 sufficiently distinct scales and orientations, given by

$$203 \quad (\phi_1, T_1) = \begin{cases} (45^\circ, 1 \text{ m}), & x \leq 1 \text{ m} \\ (135^\circ, 2 \text{ m}) & x > 1 \text{ m} \end{cases} \quad (18)$$

204 We use this case to verify the procedure for identifying locally-dominant scales.  
 205 In Case 2 (Fig. 3A), the field contains two zones with similar scales and ori-  
 206 entations, given by

$$207 \quad (\phi_2, T_2) = \begin{cases} (30^\circ, 2 \text{ m}), & x \leq 1 \text{ m} \\ (20^\circ, 1.6 \text{ m}) & x > 1 \text{ m} \end{cases} \quad (19)$$

208 We use this case to demonstrate the filtering procedure.

### 209 **Case 1: Two Zones with Distinct Periods and Orientations**

210 We use wavelet analysis to identify the scale and orientation pairs that are  
 211 dominant throughout the field  $z_1(x, y)$  (Fig. 2A). We evaluate (16) for a range

212 of scales ( $a = 0.2, 0.3, \dots, 4$  m) and orientations ( $\theta = 0^\circ, 5^\circ, \dots, 175^\circ$ ), and we  
213 calculate the wavelet variance using (10). The results (Fig. 2B) show two local  
214 maxima in the wavelet variance – one at  $(a_{\max}, \theta_{\max}) = (1 \text{ m}, 45^\circ)$ , and one at  
215  $(a_{\max}, \theta_{\max}) = (2 \text{ m}, 135^\circ)$ . These  $(a, \theta)$  pairs represent the globally-dominant  
216 scales and orientations of the field and are identical to the true scales and  
217 orientations in (18). Note that since we use  $k_o = 2\pi$  in (6), the local maxima  
218 of the wavelet variance occurs at a wavelet scale  $a_{\max}$  that is equal to the  
219 period,  $T$ , of the sinusoidal function (Torrence and Compo, 1998; Neupauer  
220 et al., 2006). For  $k_o \neq 2\pi$ ,  $a_{\max}$  is a function of  $T$ , but is not identically equal  
221 to  $T$ .

222 The local maxima of the wavelet variance identify globally-dominant scales  
223 and orientations, but they do not provide information about the locally-  
224 dominant orientations. To identify the orientation that is dominant at each  
225 location, we use (11) for each of the dominant scales identified above. Figure  
226 2C shows the dominant orientations at each position for  $a = 2$  m. This wavelet  
227 scale is equal to the true period ( $T = 2$  m) of the field for  $x > 1$  m. The dom-  
228 inant orientation for  $x > 1$  m is identified as  $135^\circ$ , which matches the true  
229 orientation of the field in this region. For  $a = 1$  m, the dominant orientations  
230 at each location are shown in Figure 2D. This wavelet scale is equal to the  
231 true period ( $T = 1$  m) of the field for  $x < 1$  m. For  $x < 1$  m, the dominant  
232 orientation of the field is identified as  $45^\circ$ , which matches the true orientation  
233 of the field in this region. The present method does not identify the location  
234 of the interface between the two zones with different orientations. This is the  
235 subject of future work.

237 We use wavelet analysis to identify the scale and orientation pairs that are  
 238 dominant throughout the field  $z_2(x, y)$  (Fig. 3A). We evaluate (16) for a range  
 239 of scales ( $a = 0.2, 0.3, \dots, 4$  m) and orientations ( $\theta = 0^\circ, 2^\circ, \dots, 178^\circ$ ), and  
 240 we calculate the wavelet variance using (10). The results (Fig. 3B) show one  
 241 local maximum at  $(a_{\max}, \theta_{\max}) = (2 \text{ m}, 30^\circ)$ , which represents the scale and  
 242 orientation of the more dominant (primary) zone of the field, i.e.,  $x < 1$  m.  
 243 This zone is more dominant because it covers a larger area of the domain,  
 244 and because it has a larger period. Although the field contains two zones with  
 245 different periods and orientations, the periods and orientations are sufficiently  
 246 similar so that the regions with large wavelet variance values overlap, and  
 247 the region containing the primary  $(a, \theta)$  pair masks the region containing the  
 248 secondary  $(a, \theta)$  pair.

249 We filter  $z_2(x, y)$  to remove the part that has a dominant orientation of  $30^\circ$ .  
 250 We accomplish this by reconstructing  $z_2(x, y)$  from (12) using only a subset of  
 251 the wavelet coefficients. To choose the subset of wavelet coefficients to use in  
 252 the reconstruction, we select a threshold value of the wavelet variance so that  
 253 the field is reconstructed using only the wavelet coefficients that correspond  
 254 to wavelet variance values that are above the threshold. We use a threshold  
 255 value of  $61.5 \text{ m}^4$ , which was chosen such that 75% of the total wavelet variance  
 256 is above the threshold. Figure 4A shows the wavelet variance of  $z_2(x, y)$  with  
 257 only values below the threshold wavelet variance remaining.

258 Using (12), we reconstruct  $z_2(x, y)$  using only the remaining wavelet coeffi-  
 259 cients (Fig. 4B). Let us denote this filtered version by  $\bar{z}_2(x, y)$ . Note that

260  $\bar{z}_2(x, y)$  approximates  $z_2(x, y)$  for  $x > 1$  m (the less dominant zone); while for  
261  $x < 1$  m (the more dominant zone),  $\bar{z}_2(x, y) \approx 0$ . Thus, the primary feature  
262 of the original field has been removed. Note that the reconstruction is not  
263 exact near the boundaries of the domain or at the boundaries between the  
264 two zones.

265 Finally, we perform wavelet analysis on  $\bar{z}_2(x, y)$  to obtain its wavelet variance.  
266 The wavelet variance (Fig. 4C) has a local maximum at  $(a_{\max}, \theta_{\max}) = (1.6$   
267  $\text{m}, 14^\circ)$ . This  $a_{\max}$  is identical to the period of  $z_2(x, y)$  for  $x > 1$  m, and  
268 this  $\theta_{\max}$  is approximately equal to the orientation of  $z_2(x, y)$  for  $x > 1$  m.  
269 Thus, the combination of filtering, reconstruction, and reanalysis allows us to  
270 identify the secondary feature of  $z_2(x, y)$ .

271 To identify the orientation that is dominant at each location, we use (11)  
272 for the two dominant scales identified above. Figure 3C shows the locally-  
273 dominant orientations for  $a = 2$  m, which is the true period for  $x < 1$  m. The  
274 wavelet analysis procedure correctly identifies the locally-dominant orientation  
275 to be  $30^\circ$  in this zone ( $x < 1$  m). Figure 3D shows the dominant orientations  
276 at each position for  $a = 1.6$  m, which is the true period for  $x > 1$  m. The  
277 wavelet analysis procedure correctly identifies the locally-dominant orientation  
278 to be  $20^\circ$  in most of this zone; however, the identified dominant orientation is  
279 slightly off near the interface between the two zones.

## 280 APPLICATION TO MASSILLON SANDSTONE DATA

281 Tidwell and Wilson (2000) collected permeability measurements (Fig. 5A) on  
282 a  $0.94 \times 0.96 \times 1.01$  m block of Massillon sandstone using a gas multisupport

283 permeameter. These measurements were taken on a square 50-by-50 grid at  
284 a spacing of  $dx = dy = 0.0127$  m. The  $0.622 \times 0.622$  m grid was centered on  
285 the block face providing a buffer of over 0.15 m between the grid and edge  
286 of the block to avoid boundary effects. The Massillon sandstone exhibits a  
287 series of subhorizontal bounding surfaces with low permeability (spacing of  
288 0.16-0.22 m; orientation of  $\theta \approx 0^\circ$ ; Tidwell and Wilson, 2000) separated by  
289 cross-stratification sets with high permeability (spacing of approximately 0.03  
290 m; orientation of  $\theta \approx -22^\circ$ ; Tidwell and Wilson, 2000).

291 Neupauer et al. (2006) performed wavelet analysis on the Massillon perme-  
292 ability data, but the wavelet analysis results only identified the bounding  
293 surfaces, which is a more dominant feature than the cross stratification sets.  
294 In this section, we use the combined wavelet analysis and filtering technique  
295 to characterize the dominant orientation of both the bounding surfaces and  
296 the cross stratification sets.

297 We performed wavelet analysis on the Massillon permeability data and cal-  
298 culated the wavelet variance using (10). We used a range of scales of  $a =$   
299  $3dx, 4dx, \dots, 0.3937$  m, and an anisotropy ratio of  $L = 0.2$ . (Powell (2004)  
300 found that  $L = 0.2$  gave the best representation of the dominant orientations  
301 for these data.) The wavelet variance (Fig. 5B) has one local maximum at  
302  $(a_{\max}, \theta_{\max}) = (0.1819 \text{ m}, 0^\circ)$ , which corresponds to the separation distance  
303 and orientation of the bounding surfaces. We would also expect a local maxi-  
304 mum at  $(a_{\max}, \theta_{\max}) \approx (0.03 \text{ m}, -22^\circ)$ , corresponding to the cross-stratification  
305 sets; however, this maximum is masked by the high wavelet variance values  
306 corresponding to the bounding surfaces. We implement the filtering method  
307 to remove the primary feature (bounding surfaces) from the field so that the  
308 secondary feature (cross-stratification sets) can be identified.

309 We filter the Massillon permeability data by reconstructing the data with the  
310 wavelet coefficients corresponding to the wavelet variance values below the  
311 threshold value ( $4.9 \times 10^{-29} \text{ m}^8$ , the value that removes 75% of the wavelet  
312 variance). The wavelet variance (Fig. 6A) from the remaining wavelet coeffi-  
313 cients shows that the original local maxima is removed. The reconstruction  
314 of the permeability data from these remaining wavelet coefficients (Fig. 6B)  
315 shows that the low-permeability bounding surfaces are removed or subdued.

316 We perform wavelet analysis on the filtered data set to identify previously-  
317 masked secondary features. The resulting wavelet variance is shown in Figure  
318 6C. The wavelet variance has a local maximum at  $(a_{\max}, \theta_{\max}) = (0.0508 \text{ m}, -20^\circ)$ .  
319 This orientation closely matches the orientation ( $-22^\circ$ ) that Tidwell and Wil-  
320 son (2000) identified in their variogram analysis. Through visual inspection of  
321 the Massillon permeability data, Tidwell and Wilson (2000) identified the sep-  
322 aration distance of cross-stratification sets to be slightly less (approximately  
323 0.03 m) than the results of our wavelet analysis.

324 For the secondary dominant scale ( $a = 0.0508 \text{ m}$ ), we use (11) to identify the  
325 locally-dominant orientations (Fig. 6D). Throughout most of the domain, the  
326 dominant orientations range from  $-30^\circ$  to  $0^\circ$ , which matches the dominant ori-  
327 entation of the cross-stratification sets. Near  $y \approx 0.7 \text{ m}$  the identified locally-  
328 dominant orientation is  $\theta \approx 0^\circ$ . This is consistent with the reconstructed data  
329 (Fig. 6B) in that region, and indicates that the filtering subdued, but did not  
330 completely remove, the bounding surface in that region. Similar behavior is  
331 seen near  $y \approx 0.55 \text{ m}$  for  $x < 0.6 \text{ m}$ . In the lower right corner of the domain,  
332 the identified locally-dominant orientations are nearly vertical, which matches  
333 both the original data (Fig. 5A) and the filtered data (Fig. 6B). Since filtering  
334 only removed wavelet coefficients that correspond to  $\theta \approx 0^\circ$ , in this region

335 where the dominant orientation is  $\theta \approx 90^\circ$ , the filtered data have the same  
336 dominant orientation as the original data.

337

## CONCLUSION

338 In this paper, we presented a wavelet analysis approach that can be used to  
339 identify dominant orientations in both stationary and non-stationary perme-  
340 ability fields. We addressed a class of non-stationarity in which the permeabil-  
341 ity field contains zones with different dominant orientations and scales, but  
342 within a given zone, the properties are stationary.

343 Wavelet analysis involves the integral transform of a data set using a wavelet  
344 as a kernel. The result of the integral transform is a set of wavelet coefficients  
345 that are used to obtain the wavelet variance. Large values of the wavelet  
346 variance occur at globally-dominant scales and orientations. If a permeability  
347 field contains two zones with distinct dominant scales and orientations, the  
348 wavelet variance will have two distinct local maxima, and the dominant scales  
349 and orientations can easily be identified. If the permeability field contains  
350 two zones with similar dominant scales and orientations, the wavelet variance  
351 may only have one local maximum, with the wavelet variance of the primary  
352 region masking the wavelet variance of the secondary region. In this paper,  
353 we developed a filtering approach that is used in conjunction with wavelet  
354 analysis to identify both primary and secondary features.

355 In the filtering approach, the original data are filtered by removing all wavelet  
356 coefficients that correspond to wavelet variance values above a chosen thresh-  
357 old. The filtered data are reconstructed from only the remaining wavelet coef-

358 ficients, thus eliminated or subduing the more dominant features of the orig-  
359 inal data. We have obtained reasonable results by using a threshold value  
360 such that 75% of the wavelet variance is removed. We illustrated the new  
361 filtering method using laboratory-collected permeability data from Massillon  
362 sandstone. The Massillon sandstone is characterized by low-permeability sub-  
363 horizontal bounding surfaces that separate low angle cross-stratification sets  
364 (Tidwell and Wilson, 2000). In prior work (Neupauer et al., 2006), wavelet  
365 analysis was used to identify dominant orientations in the Massillon sandstone  
366 permeability data; however, it was able to identify only the primary feature  
367 (bounding surfaces). With the new combined wavelet analysis and filtering  
368 approach, we identified both the primary (bounding surfaces) and secondary  
369 (cross-stratification sets) features.

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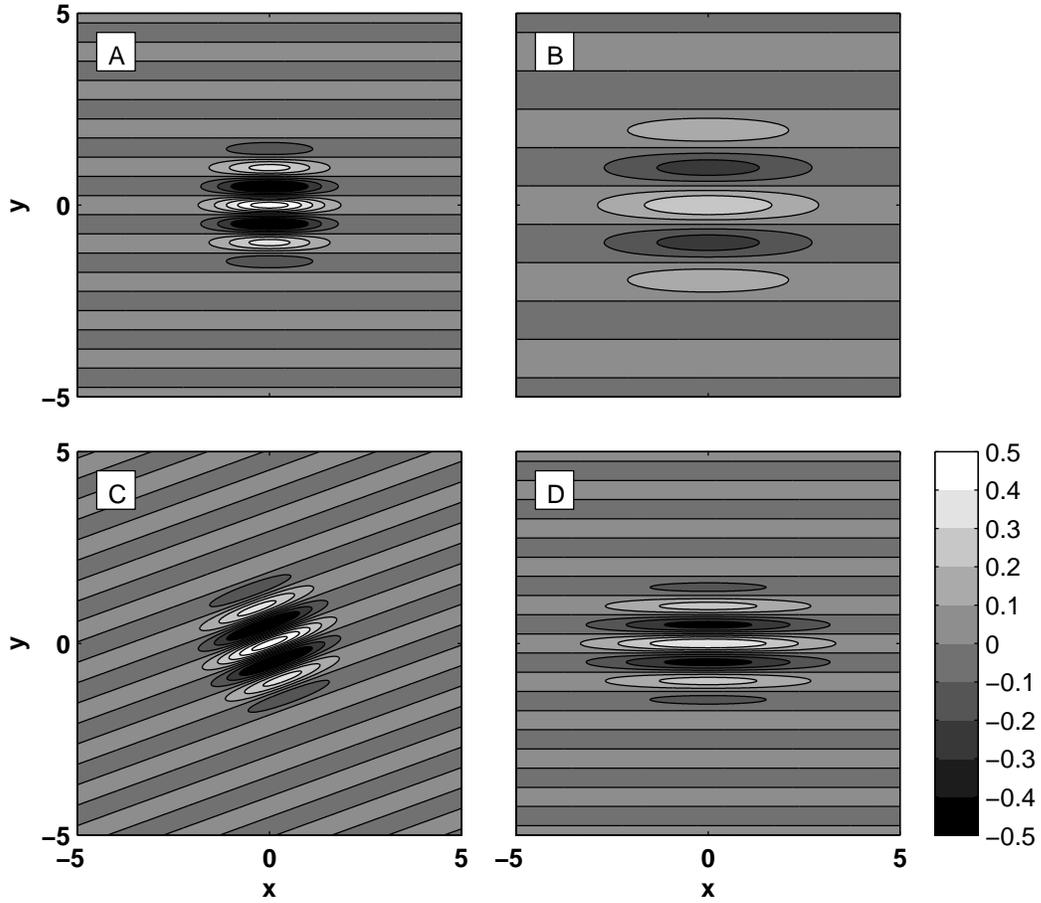


Fig. 1. Morlet wavelet (real part only). (A)  $a = 1, \mathbf{b} = (0, 0), \theta = 0^\circ, L = 1$ ;  
 (B)  $a = 2, \mathbf{b} = (0, 0), \theta = 0^\circ, L = 1$ ; (C)  $a = 1, \mathbf{b} = (0, 0), \theta = 20^\circ, L = 1$ ;  
 (D)  $a = 1, \mathbf{b} = (0, 0), \theta = 0^\circ, L = 1/2$ .

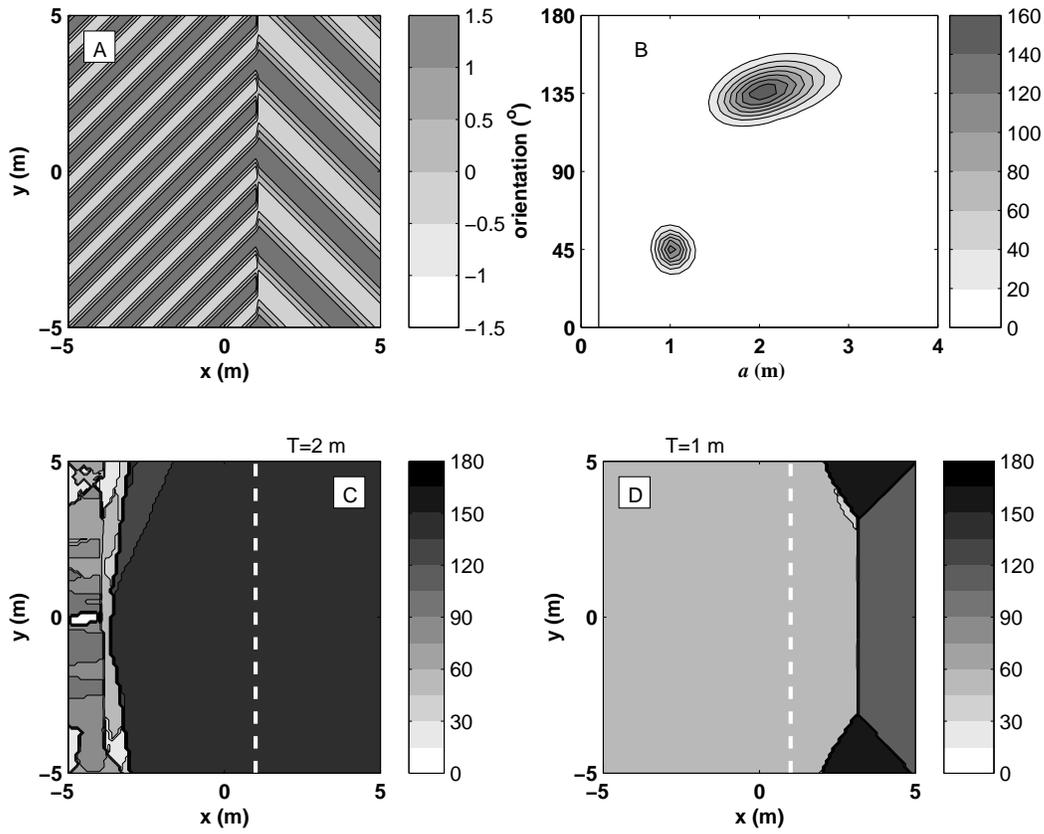


Fig. 2. Wavelet analysis of  $z_1(x, y)$ . (A)  $z_1(x, y)$ ; (B) Wavelet variance ( $\text{m}^4$ ); (C) Locally-dominant orientations (in degrees) for  $a = 2$  m; (D) Locally-dominant orientations (in degrees) for  $a = 1$  m. The dashed white line denotes the interface between the two zones. The labels above subplots C and D show the true period of the field, in the region where the true period matches the scale of the analyzing wavelet.

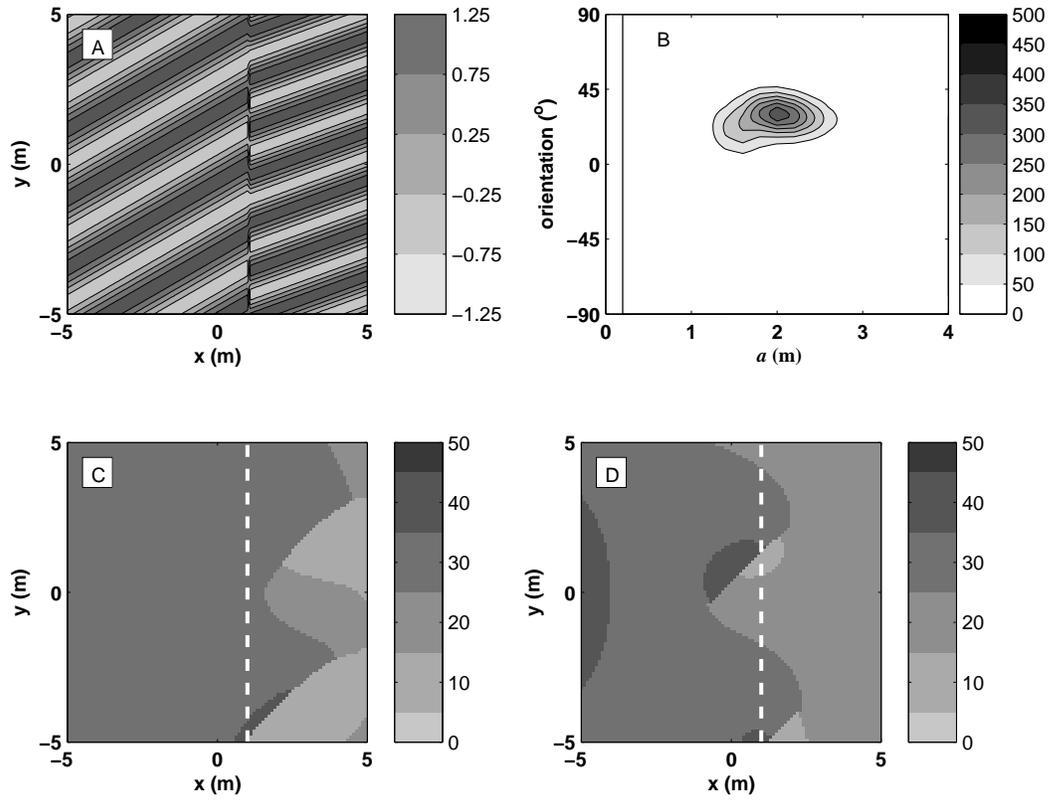


Fig. 3. Wavelet analysis of  $z_2(x, y)$ . (A)  $z_2(x, y)$ ; (B) Wavelet variance ( $\text{m}^4$ ); (C) Locally-dominant orientations (in degrees) for  $a = 2$  m; (D) Locally-dominant orientations (in degrees) for  $a = 1.6$  m. The dashed white line denotes the interface between the two zones.

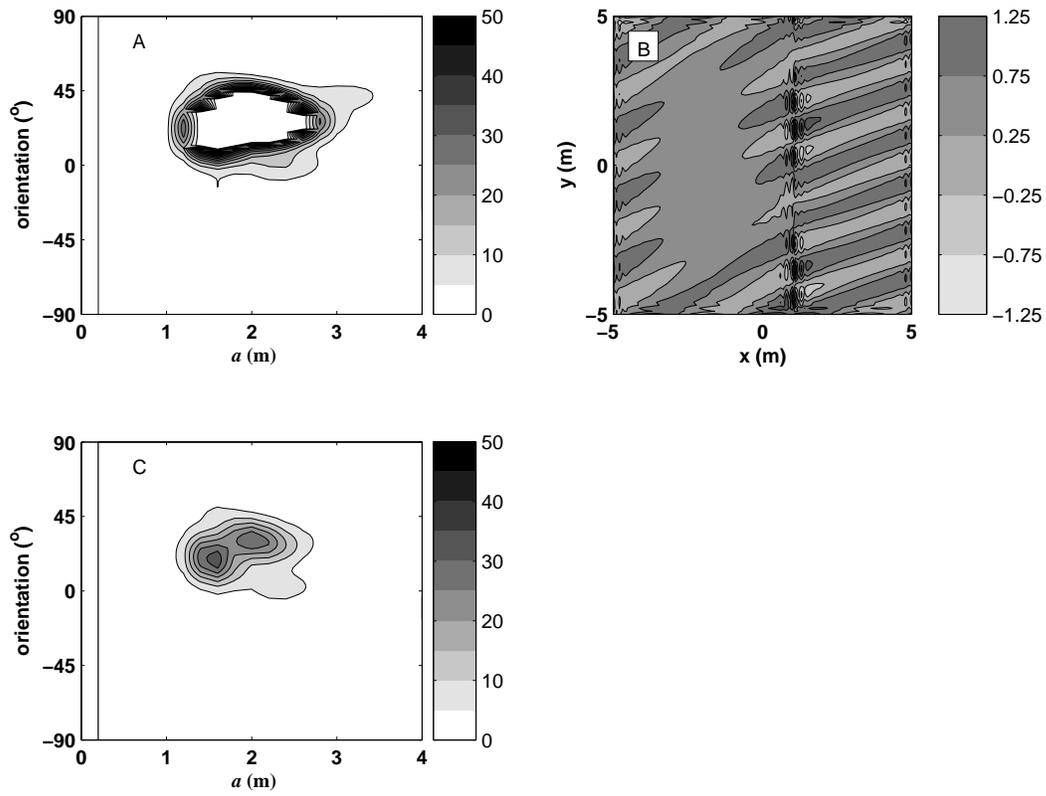


Fig. 4. Filtering and re-analysis of  $z_2(x, y)$ . (A) Wavelet variance with only values below the threshold value of  $61.5 \text{ m}^4 \text{ (m}^3\text{)}$ ; (B)  $\bar{z}_2(x, y)$ ; (C) Wavelet variance of  $\bar{z}_2(x, y) \text{ (m}^3\text{)}$ .

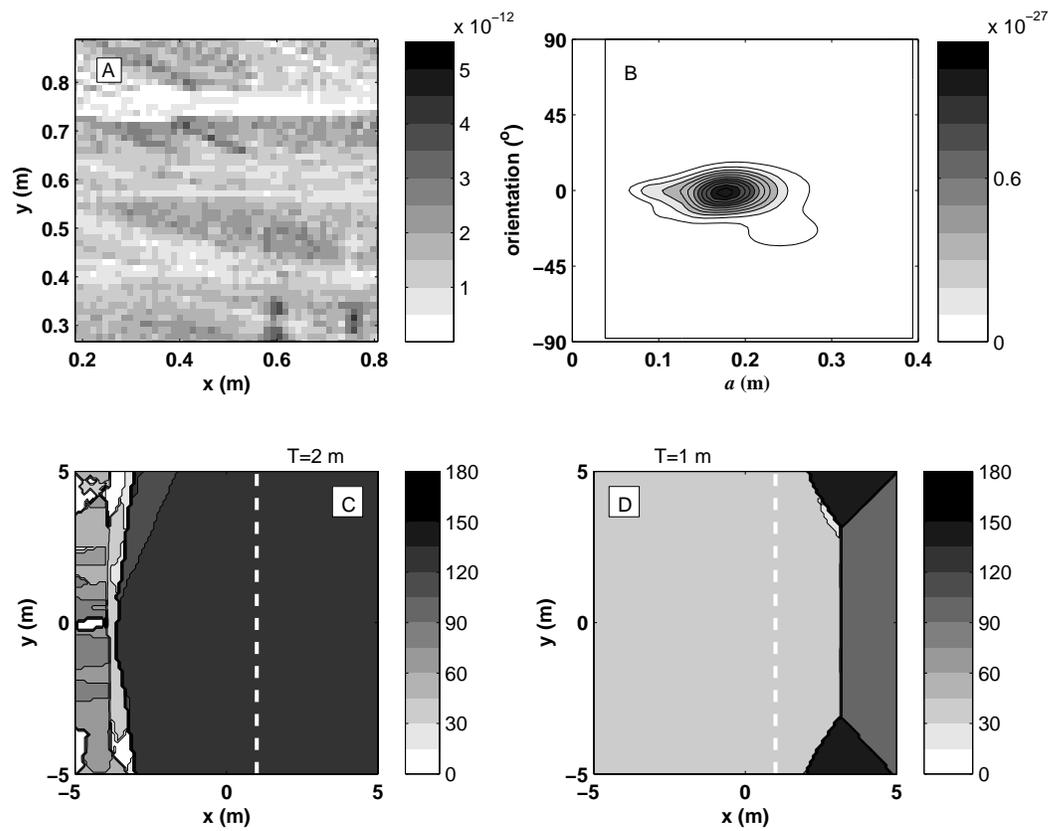


Fig. 5. Analysis of Massillon permeability data. (A) Permeability measurements ( $\text{m}^2$ ). (B) Wavelet variance. ( $\text{m}^8$ ).

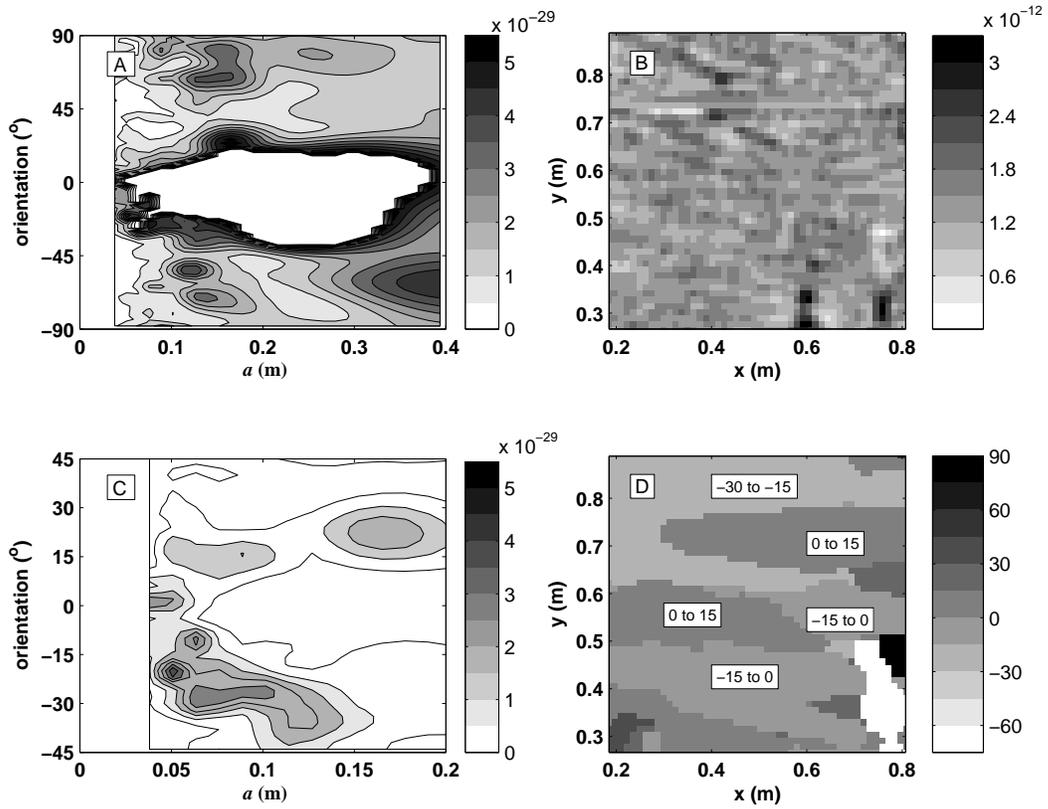


Fig. 6. Filtering, reconstruction and reanalysis of Massillon permeability data. (A) Filtered wavelet variance. (m<sup>8</sup>). (B) Reconstruction of the Massillon permeability data from the filtered wavelet coefficients (m<sup>2</sup>). (C) Wavelet variance of reconstructed Massillon permeability data (m<sup>8</sup>). (D) Dominant orientations at each location at a wavelet scale of  $a = 0.0508$  m.